

EXPLICIT AND EXACT NON-TRAVELING WAVE SOLUTIONS OF (3+1)-DIMENSIONAL GENERALIZED SHALLOW WATER EQUATION

Jianguo Liu^{1,†}, Wenhui Zhu², Li Zhou¹ and Yan He¹

Abstract In this paper, a (3+1)-dimensional generalized shallow water equation is considered. New exact solutions in forms of the hyperbolic functions and the trigonometric functions are obtained based on an extended (G'/G)-expansion method and the variable separation method, which contain traveling wave solutions and non-traveling wave solutions. The particular localized excitations and the interactions between two solitary waves for these obtained exact solutions are shown in some three-dimensional graphics.

Keywords Extended (G'/G)-expansion method, (3+1)-dimensional generalized shallow water equation, non-traveling wave solutions.

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1. Introduction

Nonlinear evolution equations (NLEEs) have been used to represent various nonlinear phenomenas in fluid dynamics, plasma physics, nonlinear optics, solid state physics, biological molecules and so on [5, 7, 9, 12, 13, 17–19, 30, 36, 37]. To understand these physical phenomenas, searching for exact solutions of NLEEs is of great important. Many methods to have been proposed [1–4, 6, 8, 10, 14, 15, 20–26, 31, 35].

Shallow water equations have applications in weather simulations, tidal waves, river and irrigation flows, tsunami prediction and so on [27]. In this work, based on the (G'/G)-expansion method and symbolic computation, we will consider the following (3+1)-dimensional generalized shallow water equation [32]

$$u_{yt} - u_{xz} - 3u_x u_{xy} - 3u_y u_{xx} + u_{xxx} = 0, \quad (1.1)$$

where $u = u(x, y, z, t)$. Eq. (1.1) describes the propagation of long water waves in oceans, estuaries, and impoundments. Tian [29] obtained the soliton-type solutions to Eq. (1.1) by using the generalized *tanh* algorithm method. Zayed [33] constructed the traveling wave solutions of Eq. (1.1) by utilizing the (G'/G)-expansion method. Tang [28] derived the Grammian and Pfaffian solutions of Eq. (1.1) by the Hirota's bilinear form. Multiple-soliton solutions were derived by Zeng [34]. Liu [11]

[†]the corresponding author. Email address: 20101059@jxutcm.edu.cn (J. Liu)

¹College of Computer, Jiangxi University of Traditional Chinese Medicine, Jiangxi 330004, China

²Institute of artificial intelligence, Nanchang Institute of Science and Technology, Jiangxi 330108, China

presented new periodic solitary wave solutions. Meng [16] studied the rational solutions of Eq. (1.1). We will discuss the non-traveling wave exact solutions by using an extended (G'/G) -expansion method, which are different from those presented in Refs. [28, 29, 33].

The organization of this paper is as follows. Section 2 proposes an extended (G'/G) -expansion method and obtains new exact solutions for the (3+1)-dimensional generalized shallow water equation. Some special soliton-structure excitations are shown by some three-dimensional graphics. Section 3 lists the discussion and summary.

2. The extended (G'/G) -expansion method and exact non-traveling wave solutions

Considering the following NLEE

$$F(u, u_x, u_y, u_z, u_t, u_{xy}, u_{xz}, u_{xt}, u_{yt}, u_{xx}, u_{tt}, \dots) = 0. \quad (2.1)$$

For finding the exact solutions of Eq. (2.1), we suppose

$$u = \sum_{i=-m}^m a_i \left(\frac{G'}{G}\right)^i, \quad (2.2)$$

where $G = G(\vartheta)$, $\vartheta = \vartheta(x, y, z, t)$, $a_i (i = -m, \dots, m)$ is unknown constant. Eq. (2.2) contains more arbitrary parameters than previous work [28, 29, 33]. The G satisfies the second-order linear ordinary differential equation

$$G'' + \lambda G' + \rho G = 0. \quad (2.3)$$

The general solutions of Eq. (2.3) are presented as follows

$$\frac{G'}{G} = \begin{cases} -\frac{\rho}{2} + \zeta_1 \frac{C_1 \cosh(\zeta_1 \vartheta) + C_2 \sinh(\zeta_1 \vartheta)}{C_1 \sinh(\zeta_1 \vartheta) + C_2 \cosh(\zeta_1 \vartheta)}, \varrho^2 - 4\rho > 0, \zeta_1 = \frac{\sqrt{\varrho^2 - 4\rho}}{2}, \\ -\frac{\rho}{2} + \zeta_2 \frac{-C_1 \sin(\zeta_2 \vartheta) + C_2 \cos(\zeta_2 \vartheta)}{C_1 \cos(\zeta_2 \vartheta) + C_2 \sin(\zeta_2 \vartheta)}, \varrho^2 - 4\rho < 0, \zeta_2 = \frac{\sqrt{-\varrho^2 + 4\rho}}{2}, \\ -\frac{\rho}{2} + \frac{C_2}{C_1 + C_2 \vartheta}, \varrho^2 - 4\rho = 0, \end{cases} \quad (2.4)$$

where ϱ , ρ , C_1 and C_2 are arbitrary constants.

Using the homogenous balance principle, we have $m = 1$. Thus Eq. (2.2) can be changed into

$$u = a_{-1} \left(\frac{G'}{G'}\right) + a_0 + a_1 \left(\frac{G'}{G}\right). \quad (2.5)$$

Supposing $\vartheta(x, y, z, t) = f(y + cz) + ax + h(t)$, substituting Eq. (2.5) into Eq. (1.1) and setting the coefficients of all terms with the same powers of $(G'/G)^k (k = -5, \dots, -2, -1, 1, 2, \dots, 5)$ to zero, we have

Case 1:

$$a_1 = 0, a_{-1} = 2a\rho, h(t) = At + B,$$

$$a_0 = cy + \frac{[(\varrho^2 - 4\rho) a^3 - 4ca + A]f}{3a^2} + W(z, t), \tag{2.6}$$

where $W(z, t)$ and $f = f(y + cz)$ are arbitrary functions, a, c, A and B are arbitrary constants.

Substituting Eq. (2.6) and the general solutions of Eq. (2.3) into Eq. (2.5), we can present new non-traveling wave solutions for Eq. (1.1).

(1): When $\varrho^2 - 4\rho > 0$, the first non-traveling wave solutions of Eq. (1.1) can be written as

$$u_1 = cx + \frac{f[(\varrho^2 - 4\rho) a^3 - 4ca + A]}{3a^2} + W(z, t) + [4a\rho C_1 \sinh[\zeta_1(B + ax + At + f)] + 4a\rho C_2 \cosh[\zeta_1(B + ax + At + f)]] / [(2\zeta_1 C_1 - \varrho C_2) \cosh[\zeta_1(B + ax + At + f)] + (2\zeta_1 C_2 - \varrho C_1) \sinh[\zeta_1(B + ax + At + f)]] \tag{2.7}$$

The physical structures of the solution (2.7) are shown in Figs.1-4.

(2): When $\varrho^2 - 4\rho < 0$, we have the second non-traveling wave solutions of Eq. (1.1)

$$u_2 = cx + \frac{f[(\varrho^2 - 4\rho) a^3 - 4ca + A]}{3a^2} + W(z, t) + 4a\rho[C_1 \cos((B + f + At + ax)\zeta_2) + C_2 \sin((B + f + At + ax)\zeta_2)] / [(2C_2\zeta_2 - \varrho C_1) \cos((B + f + At + ax)\zeta_2) - (\varrho C_2 + 2C_1\zeta_2) \sin((B + f + At + ax)\zeta_2)] \tag{2.8}$$

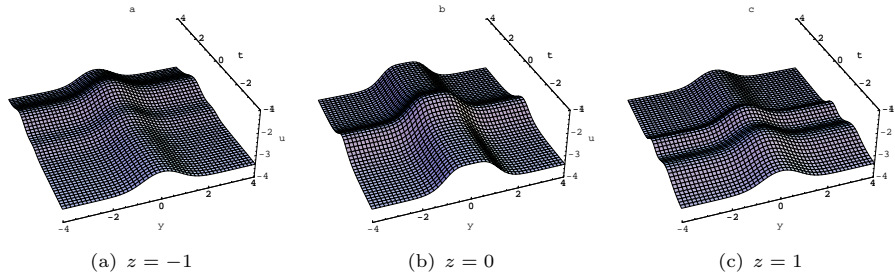


Figure 1. Solution (2.7) at $f = 1 + \text{sech}^2(y + cz)$, $W(z, t) = \text{sech}(z^2 + t^2)$, $C_1 = 1, C_2 = 2, \varrho = 3, \rho = 0.1, c = 2, x = 1, a = A = 1, B = 0$.

Case 2:

$$a_1 = -2a, a_{-1} = 2a\rho, h(t) = At + B, \varrho = 0, a_0 = cx + \frac{[(\varrho^2 - 16\rho) a^3 - 4ca + A]f}{3a^2} + W(z, t), \tag{2.9}$$

where $W(z, t)$ and $f = f(y + cz)$ are arbitrary functions, a, c, A and B are arbitrary constants.

Substituting Eq. (2.9) and the general solutions of Eq. (2.3) into Eq. (2.5), we can present another new non-traveling wave solutions for Eq. (1.1) as follows

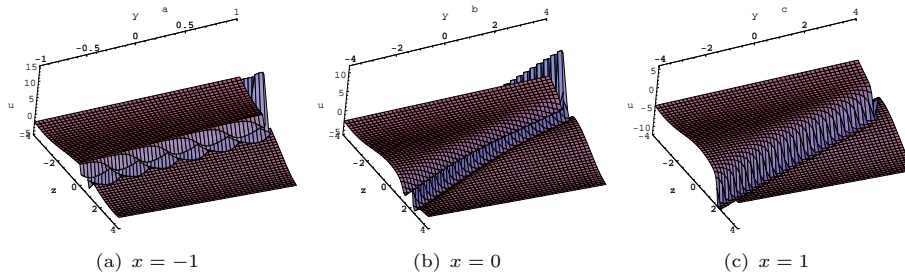


Figure 2. Solution (2.7) at $W(z, t) = \cos \sqrt{z^2 - t^2}$, $B = 0$, $C_1 = 1$, $C_2 = 2$, $f = \coth^2(y + cz) + \operatorname{csch}^2(y + cz)$, $\varrho = 3$, $\rho = 0.1$, $c = 2$, $t = 1$, $a = A = 1$.

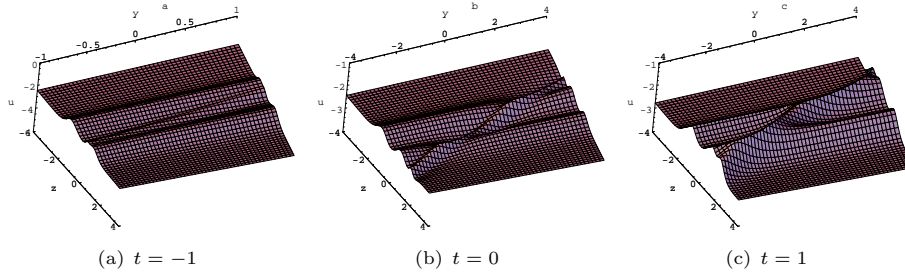


Figure 3. Solution (2.7) at $c = 2$, $f = \frac{1}{1+0.2\operatorname{sech}^2(y+cz)+0.2\tanh^2(y+cz)} - \operatorname{sech}^2(y+cz)$, $\varrho = 3$, $\rho = 0.1$, $C_2 = 2$, $a = A = 1$, $B = 0$, $W(z, t) = \tanh(z^2 + t^2) + \operatorname{sech}(t^2 - z^2)$, $x = 1$.

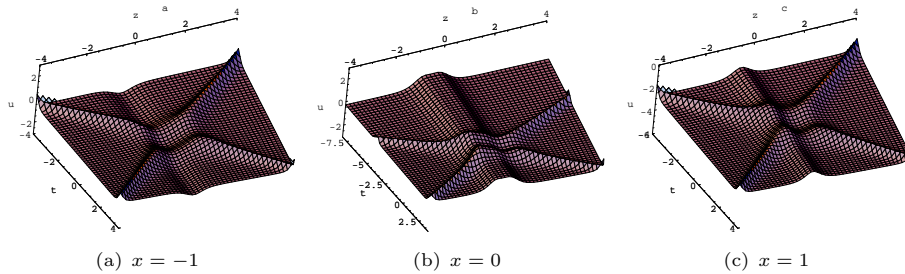


Figure 4. Solution (2.7) at $c = 2$, $f = \frac{1}{1+0.2\operatorname{sech}^2(y+cz)+0.2\tanh^2(y+cz)} - \operatorname{sech}^2(y+cz)$, $\varrho = 3$, $\rho = 0.1$, $C_2 = 2$, $a = A = 1$, $B = 0$, $W(z, t) = \tanh(z^2 + t^2) + \operatorname{sech}(t^2 - z^2)$, $y = 1$.

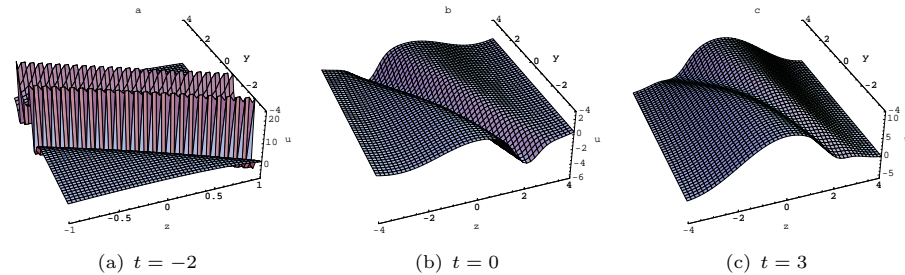


Figure 5. Ssolution (2.10) at $W(z, t) = \cos \sqrt{z^2 - t^2}$, $C_1 = 1$, $\varrho = 3$, $\rho = -1$, $c = 2$, $f = \frac{1}{\coth^2(y+cz)+\operatorname{csch}^2(y+cz)}$, $C_2 = 2$, $x = 1$, $B = 0$, $a = A = 1$.

(1): When $\rho < 0$, the third non-traveling wave solutions of Eq. (1.1) can be obtained as

$$\begin{aligned}
 u_3 = & cx + \frac{f[A - 4a(4\rho a^2 + c)]}{3a^2} + W(z, t) \\
 & - 2a\sqrt{-\rho}[C_1 \sinh[\sqrt{-\rho}(B + At + ax + f)] + C_2 \cosh[\sqrt{-\rho}(B + At + ax + f)]] \\
 & / [C_1 \cosh[\sqrt{-\rho}(B + At + ax + f)] + C_2 \sinh[\sqrt{-\rho}(B + At + ax + f)]] \\
 & - 2a\sqrt{-\rho}[C_1 \cosh[(B + f + At + ax)\sqrt{-\rho}] + C_2 \sinh[(B + f + At + ax)\sqrt{-\rho}]] \\
 & / [C_1 \sinh[(B + f + At + ax)\sqrt{-\rho}] + C_2 \cosh[(B + f + At + ax)\sqrt{-\rho}]]. \quad (2.10)
 \end{aligned}$$

The physical structures of the solution (2.10) are shown in Figs. 5-7.

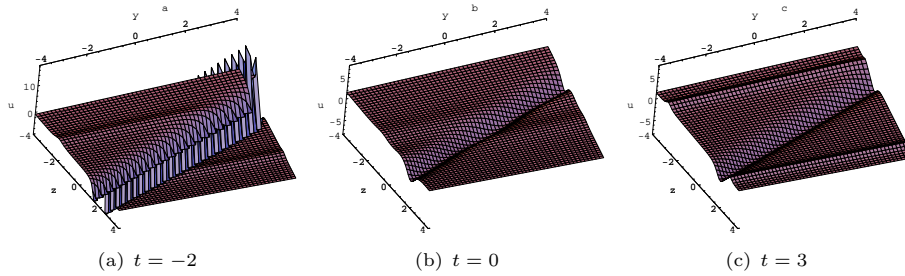


Figure 6. Solution (2.10) at $\varrho = 3$, $f = \frac{1}{1+0.2\operatorname{sech}^2(y+cz)+0.2\tanh^2(y+cz)} - \operatorname{sech}^2(y + cz)$, $\rho = -1$, $c = 2$, $C_2 = 2$, $a = A = 1$, $B = 0$, $W(z, t) = \tanh(z^2 + t^2) + \operatorname{sech}(t^2 - z^2)$, $x = 1$.

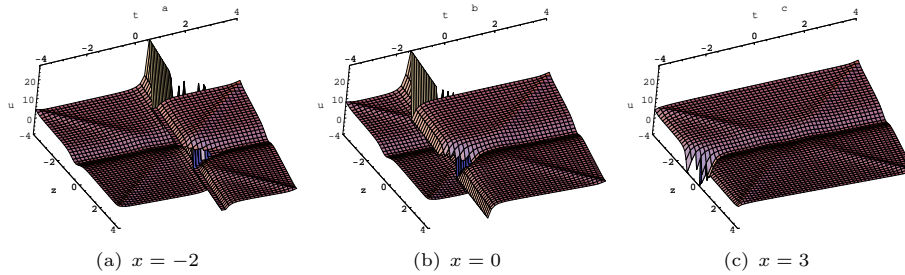


Figure 7. Solution (2.10) at $\varrho = 3$, $f = \frac{1}{1+0.2\operatorname{sech}^2(y+cz)+0.2\tanh^2(y+cz)} - \operatorname{sech}^2(y + cz)$, $\rho = 0.1$, $c = 2$, $C_2 = 2$, $a = A = 1$, $B = 0$, $W(z, t) = \tanh(z^2 + t^2) + \operatorname{sech}(t^2 - z^2)$, $y = 1$.

(2): When $\rho > 0$, the fourth non-traveling wave solutions of Eq. (1.1) can be expressed as

$$\begin{aligned}
 u_4 = & cx + \frac{f[A - 4a(4\rho a^2 + c)]}{3a^2} + W(z, t) \\
 & + 2a\sqrt{\rho}[C_1 \sin[\sqrt{\rho}(B + At + ax + f)] - C_2 \cos[\sqrt{\rho}(B + At + ax + f)]] \\
 & / [C_1 \cos[\sqrt{\rho}(B + At + ax + f)] + C_2 \sin[\sqrt{\rho}(B + At + ax + f)]] \\
 & + 2a\sqrt{\rho}[C_1 \cos[(B + f + At + ax)\sqrt{\rho}] + C_2 \sin[(B + f + At + ax)\sqrt{\rho}]] \\
 & / [-C_1 \sin[(B + f + At + ax)\sqrt{\rho}] + C_2 \cos[(B + f + At + ax)\sqrt{\rho}]]. \quad (2.11)
 \end{aligned}$$

In Figs.1-7, we research the excitation process of a special dromion soliton structure of solutions (2.7) and (2.10) for Eq. (1.1). The variation of solutions (2.7) and (2.10) with time and the interactions between two solitary waves are also described. It is obviously that other choices of $f(y + cz)$ and $W(z, t)$ in solutions (2.7) and (2.10) may form rich localized soliton structures. In other words, solutions (2.8) and (2.11) may also be employed to excite abundant soliton structures.

3. Discussion and summary

In this paper, we have obtained some new exact non-traveling wave solutions for Eq. (1.1) by using an extended (G'/G) -expansion method. Furthermore, by selecting the different values for $\vartheta(x, y, z, t)$ in the solutions (2.7) and (2.10), we can see various interesting localized soliton excitations. Also, the models proposed in this article describe important applications in physics and engineering.

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