

EXACT TRAVELLING WAVE SOLUTIONS TO THE SPACE-TIME FRACTIONAL CALOGERO-DEGASPERIS EQUATION USING DIFFERENT METHODS

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Abstract In this paper, we employed the ansatz method, the exp-function method and the $\left(\frac{G'}{G}\right)$ -expansion method for the first time to obtain the exact and traveling wave solutions of the space time fractional Calogero Degasperis equation. As a result, we obtained some soliton and traveling wave solutions for this equation by means of proposed three analytical methods and the aid of commercial software Maple. The results show that these methods are effective and powerful mathematical tool for solving nonlinear FDEs arising in mathematical physics.

Keywords Exact solution, exp-function method, $\left(\frac{G'}{G}\right)$ -expansion method, ansatz method, space time fractional Calogero-Degasperis equation.

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1. Introduction

Fractional differential equations (FDEs) are generalizations of classical differential equations of integer order. In recent years, this equations have gained considerable interest. Many significant phenomena in natural science and engineering such as fluid dynamics, biology, polymeric materials, damping law, diffusion processes, quantum, system identification, electromagnetic, mechanics, fluid flow, finance, viscoelasticity, quantum, chemistry, signal processing, control theory and so on can be modeled by this equations [39, 43, 46, 47].

Beacuae of its potential applications, scientists have devoted remarkable attempt to study the numerical and analytic solutions of nonlinear FDEs. Nevertheless, not all nonlinear FDEs are solvable. During recent years, mathematicians and physicists have developed various techniques to find solutions of nonlinear FDEs such as the variational iteration method, adomian decomposition method, the homotopy perturbation method, the fractional sub-equation method, the first integral method, the (G'/G) -expansion method, the functional variable method, the exp-function method, the modified Kudryashov method, the fractional MSE method, the ansatz method and the modified trial equation method and so on [3–10, 15, 17–31, 35, 36, 40–42, 44, 48–50, 52, 54, 55].

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There are different kinds of fractional derivative operators. The most famous one is the Caputo definition that the function should be differentiable [16]. In [37, 38], modified Riemann-Liouville derivative is proposed by Jumarie. With this kind of fractional derivative and fractional complex transform, we can convert fractional differential equations into integer-order differential equations. The order α of Jumarie’s derivative is defined by

$$D_t^\alpha u = \begin{cases} \frac{1}{\Gamma(1-\alpha)} \frac{d}{dt} \int_0^t (t-\xi)^{-\alpha} (u(\xi) - u(0)) d\xi & , 0 < \alpha < 1 \\ (f^{(n)}(t))^{(\alpha-n)}, n \leq \alpha < n+1, n \geq 1, \end{cases} \quad (1.1)$$

where $f : R \rightarrow R, t \rightarrow f(t)$ denotes a continuous (but not necessarily first-order-differentiable) function. Some useful formulas:

Property 1.

$$D_t^\alpha t^r = \frac{\Gamma(1+r)}{\Gamma(1+r-\alpha)} t^{r-\alpha}, \quad (1.2)$$

Property 2.

$$D_t^\alpha \{C_1 f(t) + C_2 g(t)\} = C_1 D_t^\alpha f(t) + C_2 D_t^\alpha g(t), \quad (1.3)$$

Property 3.

$$D_t^\alpha C = 0, \quad (1.4)$$

where C, C_1 and C_2 are constants.

In this line of thought, the structure of the article is as follows. In the next section, we briefly give the steps of the methods. In Section 3 we apply the methods to solve for the space–time fractional Calogero-Degasperis (CD) equation. Finally, conclusions are presented in last section.

2. Methods and theirs algorithms

We consider the following general (2+1) dimensional nonlinear space-time FDEs of the type

$$Q(u, D_t^\alpha u, D_x^\alpha u, D_y^\alpha u, D_t^{2\alpha} u, D_t^\alpha D_x^\alpha u, D_x^{2\alpha} u, D_x^\alpha D_y^\alpha u, D_y^{2\alpha} u, \dots) = 0, \quad (2.1)$$

where $0 < \alpha \leq 1$ and u is an unknown function, and Q is a polynomial of u and its partial fractional derivatives.

The proper traveling wave variable is [32]

$$u(x, y, t) = f(\theta), \quad (2.2)$$

$$\theta = \frac{kx^\alpha}{\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{nt^\alpha}{\Gamma(1+\alpha)}, \quad (2.3)$$

where k, m and n are nonzero constants.

By using the chain rule

$$\begin{aligned} D_t^\alpha u &= \sigma_t \frac{df}{d\theta} D_t^\alpha \theta, \\ D_x^\alpha u &= \sigma_x \frac{df}{d\theta} D_x^\alpha \theta, \\ D_y^\alpha u &= \sigma_y \frac{df}{d\theta} D_y^\alpha \theta, \end{aligned} \quad (2.4)$$

where σ_t , σ_x and σ_y are called the sigma indexes [1], and it can be take $\sigma_t = \sigma_x = \sigma_y = L$, where L is a constant.

Substituting (2.2) with (1.2) and (2.4) into (2.1), we can rewrite Eq.(2.1) in the following nonlinear ODE;

$$H(f, f', f'', f''', \dots) = 0, \quad (2.5)$$

where $\frac{df}{d\theta}$. Now we consider three different methods.

2.1. Basic idea of Ansatz method

For dark soliton solution, the starting hypothesis is in the form [11, 12, 51]

$$u(x, y, t) = A \tanh^p \theta \quad (2.6)$$

where $\theta = \frac{kx^\alpha}{\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{nt^\alpha}{\Gamma(1+\alpha)}$ and A , k , c are nonzero constants. From the ansatz given above with expression (2.6), it is possible to obtain necessary derivatives. Then, the obtained derivatives are substituted in the Eq.(2.5) and we collect all terms with the same order of necessary terms. Then by equating each coefficient of the resulting polynomial to zero, we obtain a set of algebraic equations for; A , k and c . Finally solving the system of equations we can get exact solution of Eq.(2.1).

2.2. Basic idea of Exp-function method

This method is based on the assumption that traveling wave solutions can be expressed in the following form which was developed by He and Wu [33]

$$f(\theta) = \frac{\sum_{n=-c}^d a_n \exp[n\theta]}{\sum_{m=-p}^q b_m \exp[m\theta]} \quad (2.7)$$

where p, q, c and d are positive integers, a_n and b_m are unknown constants. Also this expression can write in the following equivalent form.

$$f(\theta) = \frac{a_{-c} \exp[-c\theta] + \dots + a_d \exp[d\theta]}{b_{-p} \exp[-p\theta] + \dots + b_q \exp[q\theta]}. \quad (2.8)$$

To determine the value of c and p , we balance the linear term of lowest order of equation Eq.(2.5) with the lowest order nonlinear term. Similarly, to determine the value of d and q , we balance the linear term of highest order of Eq.(2.5) with highest order nonlinear term [2, 13, 34, 56].

2.3. Basic idea of $\left(\frac{G'}{G}\right)$ -expansion method

Step 1: According to the this method which was developed by Wang [53], we look for its solution $f(\theta)$ in the polynomial form

$$f(\theta) = \sum_{i=0}^z a_i \left(\frac{G'}{G}\right)^i, \quad a_z \neq 0, \quad (2.9)$$

where a_i are constants, while $G(\theta)$ is the solution of the auxiliary linear second order ODE

$$G''(\theta) + \lambda G'(\theta) + \mu G(\theta) = 0, \quad (2.10)$$

with λ and μ are being constants.

Step 2: z is a positive integer which is determined by the homogeneous balancing method in Eq.(2.5).

Step 3: By substituting Eqs.(2.9) and (2.10) into Eq.(2.5) with the value of z obtained in Step 2, and collecting all terms with the same order of $\left(\frac{G'}{G}\right)$ together. Then setting each coefficient to zero, we obtained a set of algebraic equations for a_i, k, m and n .

Step 4: Solve the system of algebraic equations obtained in step 3 for a_i ($i = 0, 1, 2, \dots, z$), k, m and n by use of Maple. Then we substitute a_i ($i = 0, 1, 2, \dots, z$), k, m, n and the solutions of Eq.(2.10) into Eq.(2.9), we can obtain a series of fundamental solutions of Eq.(2.1) [14].

3. New solutions of the space–time fractional Calogero–Degasperis (CD) equation

We consider the space–time fractional Calogero–Degasperis (CD) equation [45]

$$D_t^\alpha D_x^\alpha u - 4D_x^\alpha u D_x^{2\alpha} u - 2D_y^\alpha u D_x^{2\alpha} u + D_y^\alpha D_x^{3\alpha} u = 0, \quad (3.1)$$

where $0 < \alpha \leq 1$.

By using the transformations (2.2) and (2.3) with (2.4), then by once integrating, Eq.(3.1) reduced into following ODE

$$-nkf' - (2k^3L + mk^2L)(f')^2 + mk^3L^2f''' = 0, \quad (3.2)$$

where “ f' ” = $\frac{df}{d\theta}$ and constant of integration is taken to be zero.

3.1. Exact solution by ansatz method

From the ansatz given above Eq.(2.6), we get necessary derivatives and these are substituted in the Eq.(3.1) we get algebraic equation. By use this equation, in order to reduce the number of coefficients of the powers of $\tanh \theta$, we determine the balance value of p . Equating $2(p+1) = p+3$ yields $p = 1$. Using this value of $p = 1$, this algebraic equation reduces to

$$\begin{aligned} & (8A^2n^2k - 4A^2mkn + 24Ak^3m) \tanh^5 \\ & + (8A^2mkn - 40Ak^3m - 16A^2n^2k - 2Akn) \tanh^3 \\ & + (16Ak^3m + 2Akn - 4A^2mkn + 8A^2n^2k) \tanh = 0. \end{aligned} \quad (3.3)$$

So, we obtain a system of algebraic equations

$$\begin{aligned} 8A^2n^2k - 4A^2mkn + 24AL^2k^3m &= 0, \\ 8A^2mkn - 40AL^2k^3m - 16A^2n^2k - 2Akn &= 0, \\ 16AL^2k^3m + 2Akn - 4A^2mkn + 8A^2n^2k &= 0. \end{aligned} \quad (3.4)$$

Solving this system for A and n gives

$$A = \frac{6L^2 k^2 m}{n(m - 2n)}, \quad (3.5)$$

$$n = 4L^2 k^2 m. \quad (3.6)$$

Equation (3.5) prompts the constraint $m \neq 2n$. Thus finally, the dark soliton solution for Eq.(3.1) is given by

$$u(x, y, t) = A \tanh \left(\frac{kx^\alpha}{\Gamma(1 + \alpha)} + \frac{my^\alpha}{\Gamma(1 + \alpha)} - \frac{nt^\alpha}{\Gamma(1 + \alpha)} \right), \quad (3.7)$$

where A is given in Eq.(3.5) and n is given in Eq.(3.6).

3.2. Exact solutions by Exp-function method

By balancing the order of f''' and $(f')^2$ in Eq.(3.2), we obtain

$$f''' = \frac{c_1 \exp[-(c + 7p)\theta] + \dots}{c_2 \exp[-8p\theta] + \dots}, \quad (3.8)$$

and

$$(f')^2 = \frac{c_3 \exp[-(2c + 2p)\theta] + \dots}{c_4 \exp[-4p\theta] + \dots}, \quad (3.9)$$

where c_i are determined coefficients only for simplicity. Balancing lowest order of exp-function in Eqs.(3.8) and (3.9) we get

$$-(c + 7p) = -(2c + 6p), \quad (3.10)$$

and

$$p = c. \quad (3.11)$$

Similar process we balance the linear term of highest order in Eq.(3.2), we obtain

$$q = d. \quad (3.12)$$

The simplest choice for c , p , q and d is $p = c = 1$ and $q = d = 1$. According to $p = c = 1$ and $q = d = 1$, Eq.(2.8) becomes

$$f(\theta) = \frac{a_1 \exp(\theta) + a_0 + a_{-1} \exp(-\theta)}{b_1 \exp(\theta) + b_0 + b_{-1} \exp(-\theta)}. \quad (3.13)$$

Substituting Eq.(3.13) into Eq.(3.2), and solving this system of algebraic equations by the help of Maple, we get the following results

Case 1:

$$\begin{aligned} a_1 &= 0, & a_0 &= a_0, & a_{-1} &= a_{-1}, \\ b_1 &= 0, & b_0 &= 0, & b_{-1} &= b_{-1}, \\ k &= -\frac{m}{2}, & m &= m, & n &= \frac{m^3 L^2}{4}. \end{aligned} \quad (3.14)$$

Substituting these results into (3.13), we get the following exact solution

$$u_1(x, y, t) = \frac{a_0 + a_{-1} \exp\left(-\left(\frac{-mx^\alpha}{2\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{m^3 L^2 t^\alpha}{4\Gamma(1+\alpha)}\right)\right)}{b_{-1} \exp\left(-\left(\frac{-mx^\alpha}{2\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{m^3 L^2 t^\alpha}{4\Gamma(1+\alpha)}\right)\right)}. \quad (3.15)$$

Case 2:

$$\begin{aligned} a_1 &= a_1, & a_0 &= a_0, & a_{-1} &= a_{-1}, \\ b_1 &= 0, & b_0 &= b_0, & b_{-1} &= 0, \\ k &= -\frac{m}{2}, & m &= m, & n &= \frac{m^3 L^2}{4}. \end{aligned} \quad (3.16)$$

Substituting these results into (3.13), we obtain

$$u_2(x, y, t) = \frac{a_1 \exp\left(\frac{-mx^\alpha}{2\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{m^3 L^2 t^\alpha}{4\Gamma(1+\alpha)}\right) + a_0 + a_{-1} \exp\left(-\left(\frac{-mx^\alpha}{2\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{m^3 L^2 t^\alpha}{4\Gamma(1+\alpha)}\right)\right)}{b_0}. \quad (3.17)$$

Case 3:

$$\begin{aligned} a_1 &= a_1, & a_0 &= 0, & a_{-1} &= a_{-1}, \\ b_1 &= 0, & b_0 &= 0, & b_{-1} &= b_{-1}, \\ k &= k, & m &= -2k, & n &= -8k^3 L^2. \end{aligned} \quad (3.18)$$

Substituting these results into (3.13), we get the following exact solution

$$u_3(x, y, t) = \frac{a_1 \exp\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} - \frac{2ky^\alpha}{\Gamma(1+\alpha)} + \frac{8k^3 L^2 t^\alpha}{\Gamma(1+\alpha)}\right) + a_{-1} \exp\left(-\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} - \frac{2ky^\alpha}{\Gamma(1+\alpha)} + \frac{8k^3 L^2 t^\alpha}{\Gamma(1+\alpha)}\right)\right)}{b_{-1} \exp\left(-\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} - \frac{2ky^\alpha}{\Gamma(1+\alpha)} + \frac{8k^3 L^2 t^\alpha}{\Gamma(1+\alpha)}\right)\right)}. \quad (3.19)$$

Case 4:

$$\begin{aligned} a_1 &= 0, & a_0 &= -\frac{b_0(6kLmb_{-1} - (2k+m)a_{-1})}{b_{-1}(2k+m)}, & a_{-1} &= a_{-1}, \\ b_1 &= 0, & b_0 &= b_0, & b_{-1} &= b_{-1}, \\ k &= k & m &= m, & n &= mk^2 L^2, \end{aligned} \quad (3.20)$$

where a_{-1} , b_0 and b_{-1} are free parameters. When we substitute these results into (3.13), we get the following exact solution

$$u_4(x, y, t) = \frac{-\frac{b_0(6kLmb_{-1} - (2k+m)a_{-1})}{b_{-1}(2k+m)} + a_{-1} \exp\left(-\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{mk^2 L^2 t^\alpha}{\Gamma(1+\alpha)}\right)\right)}{b_0 + b_{-1} \exp\left(-\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{mk^2 L^2 t^\alpha}{\Gamma(1+\alpha)}\right)\right)}. \quad (3.21)$$

Case 5:

$$\begin{aligned} a_1 &= a_1, & a_0 &= \frac{a_1 b_{-1}^2 + a_{-1} b_0^2}{b_{-1} b_0}, & a_{-1} &= a_{-1}, \\ b_1 &= 0, & b_0 &= b_0, & b_{-1} &= b_{-1}, \\ k &= k & m &= -2k, & n &= -2k^3 L^2, \end{aligned} \quad (3.22)$$

where a_{-1} , b_0 and b_{-1} are free parameters. When we substitute these results into (3.13), we have

$$u_5(x, y, t) = \frac{a_1 \exp\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} - \frac{2ky^\alpha}{\Gamma(1+\alpha)} + \frac{2k^3L^2t^\alpha}{\Gamma(1+\alpha)}\right) + \frac{a_1b_{-1}^2 + a_{-1}b_0^2}{b_{-1}b_0} + a_{-1} \exp\left(-\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} - \frac{2ky^\alpha}{\Gamma(1+\alpha)} + \frac{2k^3L^2t^\alpha}{\Gamma(1+\alpha)}\right)\right)}{b_0 + b_{-1} \exp\left(-\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} - \frac{2ky^\alpha}{\Gamma(1+\alpha)} + \frac{2k^3L^2t^\alpha}{\Gamma(1+\alpha)}\right)\right)}. \quad (3.23)$$

If we take $a_0 = 1$, $a_{-1} = 1$ and $b_{-1} = 1$, solution (3.15) becomes

$$u_1(x, y, t) = \frac{1 + \cosh\left(\frac{-mx^\alpha}{2\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{m^3L^2t^\alpha}{4\Gamma(1+\alpha)}\right) - \sinh\left(\frac{-mx^\alpha}{2\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{m^3L^2t^\alpha}{4\Gamma(1+\alpha)}\right)}{\cosh\left(\frac{-mx^\alpha}{2\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{m^3L^2t^\alpha}{4\Gamma(1+\alpha)}\right) - \sinh\left(\frac{-mx^\alpha}{2\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{m^3L^2t^\alpha}{4\Gamma(1+\alpha)}\right)}. \quad (3.24)$$

Similarly; if we take $a_0 = 2$, $a_1 = 1$, $a_{-1} = 1$ and $b_0 = 2$, solution (3.17) becomes

$$u_2(x, y, t) = 2 \cosh^2\left(-\frac{mx^\alpha}{4\Gamma(1+\alpha)} + \frac{my^\alpha}{2\Gamma(1+\alpha)} - \frac{m^3L^2t^\alpha}{8\Gamma(1+\alpha)}\right). \quad (3.25)$$

When we take $a_1 = 1$, $a_{-1} = 1$ and $b_{-1} = 1$, solution (3.19) becomes

$$u_3(x, y, t) = 1 + \cosh\left(\frac{2kx^\alpha}{\Gamma(1+\alpha)} - \frac{4ky^\alpha}{\Gamma(1+\alpha)} + \frac{16k^3L^2t^\alpha}{\Gamma(1+\alpha)}\right) + \sinh\left(\frac{2kx^\alpha}{\Gamma(1+\alpha)} - \frac{4ky^\alpha}{\Gamma(1+\alpha)} + \frac{16k^3L^2t^\alpha}{\Gamma(1+\alpha)}\right). \quad (3.26)$$

Also if we take $a_{-1} = 1$, $b_{-1} = 1$ and $6kLm = 2k + m$, solution (3.21) becomes

$$u_4(x, y, t) = \frac{\cosh\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{mk^2L^2t^\alpha}{\Gamma(1+\alpha)}\right) - \sinh\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{mk^2L^2t^\alpha}{\Gamma(1+\alpha)}\right)}{1 + \cosh\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{mk^2L^2t^\alpha}{\Gamma(1+\alpha)}\right) - \sinh\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{mk^2L^2t^\alpha}{\Gamma(1+\alpha)}\right)}. \quad (3.27)$$

Finally, when we take $a_1 = 1$, $a_{-1} = -1$, $b_0 = 1$ and $b_{-1} = 1$, solution (3.23) becomes

$$u_5(x, y, t) = 1 + \cosh\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} - \frac{2ky^\alpha}{\Gamma(1+\alpha)} + \frac{2k^3L^2t^\alpha}{\Gamma(1+\alpha)}\right) - \sinh\left(\frac{kx^\alpha}{\Gamma(1+\alpha)} - \frac{2ky^\alpha}{\Gamma(1+\alpha)} + \frac{2k^3L^2t^\alpha}{\Gamma(1+\alpha)}\right). \quad (3.28)$$

3.3. Exact solution by $\left(\frac{G'}{G}\right)$ -expansion method

According to homogeneous balancing method, with balancing the f''' and $(f')^2$, we get $2z + 2 = 3 + z$, hence $z = 1$. Then we suppose that Eq.(3.2) has the following formal solution

$$f(\theta) = a_0 + a_1 \left(\frac{G'}{G}\right), \quad a_1 \neq 0. \quad (3.29)$$

By using Eq.(3.29) and Eq.(2.10) we have necessary derivaties and substituting them into Eq.(3.2), collecting the coefficients of $\left(\frac{G'}{G}\right)^i$ ($i = 0, \dots, 4$) and setting it to zero we derive a set of algebraic equations and solving this system by Maple we get

$$\begin{aligned} a_0 &= a_0, & a_1 &= -\frac{6mkL}{m+2k}, \\ k &= k, & m &= m, \\ n &= (\lambda^2 - 4\mu) mk^2L^2, \end{aligned} \quad (3.30)$$

λ and μ are arbitrary constants. By using (3.30) with expression (3.2) can be written as

$$f(\theta) = a_0 - \frac{6mkL}{m+2k} \left(\frac{G'}{G} \right). \tag{3.31}$$

When $\lambda^2 - 4\mu > 0$, substituting the general solution of (2.10) into (3.31), we obtain the following traveling wave solution of space–time fractional CD equation

$$f_1(\theta) = a_0 + \frac{3mkL\lambda}{m+2k} - \frac{3mkL\sqrt{\lambda^2 - 4\mu}}{m+2k} \left(\frac{C_1 \sinh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\theta + C_2 \cosh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\theta}{C_1 \cosh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\theta + C_2 \sinh \frac{1}{2}\sqrt{\lambda^2 - 4\mu}\theta} \right), \tag{3.32}$$

where $\theta = \frac{kx^\alpha}{\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{(\lambda^2 - 4\mu)mk^2L^2t^\alpha}{\Gamma(1+\alpha)}$.

Similarly when $\lambda^2 - 4\mu < 0$, we get

$$f_2(\theta) = a_0 + \frac{3mkL\lambda}{m+2k} - \frac{3mkL\sqrt{4\mu - \lambda^2}}{m+2k} \left(\frac{-C_1 \sin \frac{1}{2}\sqrt{4\mu - \lambda^2}\theta + C_2 \cos \frac{1}{2}\sqrt{4\mu - \lambda^2}\theta}{C_1 \cos \frac{1}{2}\sqrt{4\mu - \lambda^2}\theta + C_2 \sin \frac{1}{2}\sqrt{4\mu - \lambda^2}\theta} \right), \tag{3.33}$$

where $\theta = \frac{kx^\alpha}{\Gamma(1+\alpha)} + \frac{my^\alpha}{\Gamma(1+\alpha)} - \frac{(\lambda^2 - 4\mu)mk^2L^2t^\alpha}{\Gamma(1+\alpha)}$.

In particular, if $C_1 \neq 0, C_2 = 0, \lambda > 0, \mu = 0, a_0 = 0$ then f_1 and f_2 become

$$u(x, y, t) = \frac{3mkL\lambda}{m+2k} \left\{ 1 - \tanh \left(\frac{\lambda kx^\alpha}{2\Gamma(1+\alpha)} + \frac{\lambda my^\alpha}{2\Gamma(1+\alpha)} - \frac{\lambda^3 mk^2 L^2 t^\alpha}{2\Gamma(1+\alpha)} \right) \right\}. \tag{3.34}$$

4. Conclusion

In this paper, based on the the ansatz method, the exp-function method and the (G'/G) –expansion method, the space time fractional CD equation is solved exactly. The obtained exact solutions are either hyperbolic function solutions or turned into hyperbolic function solutions when suitable parameters are chosen. Comparing our results to the Mohyud-Din’s results [45] it can be seen that these results are new. Moreover, when the established solutions are compared with each other, it can be seen that solutions (3.7), (3.24), (3.25), (3.26), (3.27), (3.28) and (3.34) are different and never been obtained. When we choose $\lambda = 2, a_0 = -\frac{3mkL\lambda}{m+2k}$ and $L = -\frac{m+2k}{4mk(m-2n)}$, solution (3.34) can be convert solution (3.7). Also, we can compare methods. While the ansatz method and the (G'/G) –expansion method only give one solution, the exp-function method gives different and variety of travelling wave solutions. The performance of these methods are found to be reliable, effective, very powerful and convenient for solving nonlinear FDEs.

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